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NAVORD REPORT

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THE ATTENUATION OF SHOCK IN LUCITE (U)

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U. S. NAVAL ORDNANCE LABORATORY WHITE OAK, MARYLAND



THE ATTENUATION OF SHOCK IN LUCITE

Ву

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ABSTRACT: The attenuation of the velocity of a shock wave was measured in Lucite under conditions similar to those of the shock sensitivity test. Two systems, one based upon the reaction of pressure probes to the pressure pulse of the shock wave and the other a smear camera, were used to record the events. The reliability of the pressure probe in recording the events was comparable to the smear camera record of the shock for the first three inches of Lucite after which the response lagged behind the camera record. With the aid of the smear camera, additional data were calculated for Lucite in the low pressure region (4-5 kbar) by measuring the shock velocity in Lucite and water. These data were used to extend the equation of state for Lucite to the region applicable to this investigation.

The shock pressure in Lucite was calculated as a function of the Lucite length from the velocity obtained experimentally and the equation of state for Lucite. This was compared to the length of the gap in the shock sensitivity tests to obtain an approximate value of the pressure required for the initiation to detonation of various explosives.

An investigation of the effect of varying donor length on the shock transmitted into Lucite has been initiated.

PUBLISHED AUGUST 1960

CHEMISTRY RESEARCH DEPARTMENT U. S. Naval Ordnance Laboratory Silver Spring, Maryland Gap tests for explosives, in which sensitivity of an explosive is measured by interposing a gap of some inert material between a high explosive donor and the explosive under test, have been used for a number of years. The mechanism of initiation by shock and the meaning of gap sensitivity in relation to other sensitivity tests and to handling experience has now become of considerable interest. In the case of composite propellants, gap tests seem to reflect handling hazards of finished grains more accurately than the impact test. For this reason calibration of the gap test is of importance to knowledge of propellant sensitivity.

This research covers part of a program to calibrate the card gap test in terms of basic parameters; in this case, in terms of the minimum pressure required to initiate detonation under the conditions of the experiment. This represents an important advance in deducing the energy input and energy flux relations which are believed to be the basic information required to interpret gap sensitivities.

This research was supported by Task NOL-323, Polaris Sensitivity.

W. D. COLEMAN Captain, USN Commander

Albert Lighthody

ALBERT LIGHTBODY

By direction

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THE ATTENUATION OF SHOCK IN LUCITE

I. INTRODUCTION

The present investigation was made to interpret the shock sensitivity test results in terms of shock pressure rather than the gap thickness required to initiate the explosive. The present work was carried out on Lucite rods, since it was determined cellulose acetate (used in forming the gap) and Lucite were similar shock attenuators. The investigation consisted of extending the equation of state data of Lucite to the lower pressures in the gap and of using the data obtained to relate pressure and gap thickness for the conditions under which the gap tests are made.

The equation of state data were obtained by initiating a shock with two cylindrical tetryl pellets (each 2 inches dia. x l inch thick) and measuring the shock velocity as a function of distance in Lucite rods and in water (the equation of state of which is known) as it progresses from the Lucite to the water. Using the customary approximation at the Lucite-water interface, the pressure and particle velocity in Lucite before the interface may be obtained.

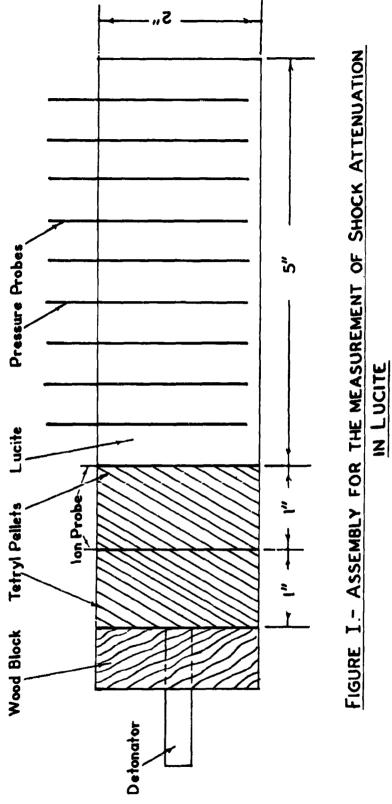
II. EXPERIMENTAL METHODS

The attenuation of the shock velocity in Lucite was determined by two different experimental techniques. One was based upon recording the passage of a pressure pulse by an electronic system, and the other used high speed photography to follow the shock front. This latter technique was used to obtain additional data to determine a more accurate curve for the equation of state of Lucite.

A. Electronic System Used to Measure Shock Velocity

Figure I is a schematic drawing of the experimental assembly used to measure the attenuation of a shock wave in a Lucite rod. A donor, consisting of a tetryl pellet or a series of tetryl pellets, was initiated by a Seismo* detonator. The detonation wave developed in the tetryl becomes a shock wave in the Lucite rod. The progress of the shock wave was followed by a series of pressure probes carefully placed in the assembly. Basically the probes act as switches which in one case were shorted by an ionization front (i.e. as an ionization probe used to trigger the measurements) and otherwise by a pressure pulse (i.e. acting as a pressure probe). The pressure pulse impinged

^{*} Detonators were obtained from Olin Mathieson.

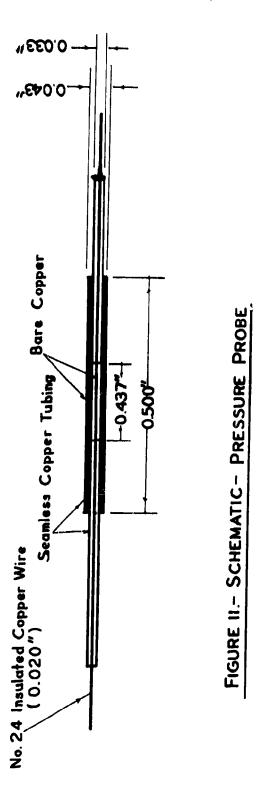


on a copper tube 0.033 inches away from a copper wire. When these two made contact, a circuit was closed and an impulse was transmitted to an oscilloscope (Tektronic No. 535). A Polaroid camera was used to make permanent records of the oscillograph tracings. Figure II is a schematic drawing of the pressure probe.

A series of holes 0.05; inches in diameter were carefully made at specified intervals in a Lucite rod. The pressure probes were inserted and the necessary leads were soldered to the probes. The tetryl pellets were securely taped to the Lucite rod. An ionization probe was inserted between the last two tetryl pellets, and at the tetryl-Lucite interface. The entire ensemble was placed in the bombproof chamber where the leads were connected to the oscilloscope and the detonator put in place. Meanwhile a series of calibrated time marks was obtained on the oscilloscope by using a Tektronic No. 181 Time-Mark Generator. The time scale was recorded on the film just prior to the experiment.

The oscilloscope was triggered by the ionization probe placed between the last two tetryl pellets. By beginning the oscilloscope sweep prior to the arrival of the detonation at the Lucite-tetryl interface, a much more definitive and precise measurement was obtained of the time of arrival of the shock in the Lucite. (In the instance when one pellet was used, the ionization probe placed at the tetryl-Lucite interface was used as a trigger.) The arrival of the reactive shock at the tetryl-Lucite interface was recorded by the second ionization probe. The further progress of the shock wave down the Lucite rod was followed by the pressure probes. A more comprehensive discussion of the pressure probe and the electronic system used is given elsewhere (1.4).

The system of most interest was the one containing a donor made up of two tetryl pellets, since these are the conditions in the shock sensitivity test. Cellulose acetate cards, 0.01 inches thick by 2 inches in diameter, are used to build gaps less than one-half inch thick. For larger gaps, Lucite discs, one-half inch and 1 inch thick are used with the cellulose acetate cards to build the required gap. A number of charges was prepared in the exact manner used for the shock sensitivity tests and ionization probes were placed at designated positions in the Lucite-cellulose acetate gap (Figure III). gap was prepared by stacking the cards and discs in units one to two inches high. Each unit was compressed to form a compact pile and a hole was drilled in it for a pressure probe. The attenuation of the shock velocity was measured at 0.5, 1.0 and 1.5 inches and compares with the shock velocity measured in the Lucite rod.



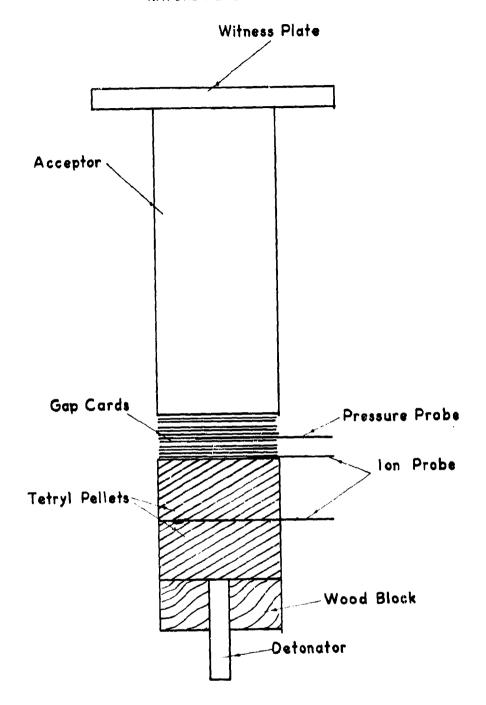


FIGURE III.- ASSEMBLY- SHOCK ATTENUATION IN THE GAP

B. High Speed Photography

The objects of this experiment were three-fold:

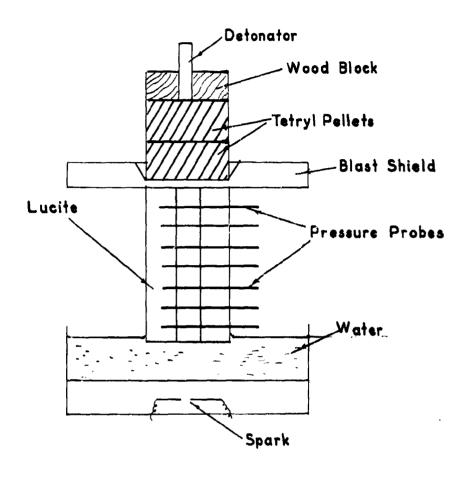
- 1) to measure the attenuation of the shock wave in Lucite by an alternate method;
- 2) to determine the reliability of the measurements made by the pressure probes; and
- 3) to obtain data which will define more precisely the equation of state of Lucite for the lower shock pressures.

Figure IV is a schematic drawing which shows the arrangement of the various components. The Lucite rod was machined from a bar 2 inches x = 1/4 inches in cross section to a rod approximately 2 1/16 inches in diameter with two parallel flat surfaces 2 inches apart and 5/8 inches wide. These parallel flats eliminated distortion of the light by the curved surfaces as the light passed through the Lucite rod. Pressure probes were inserted at designated points in the usual manner. The rod was supported vertically with its end submerged approximately 1/4 inch below the surface of the water contained in a small trough. A Lucite blast shield of known thickness was placed on top of the Lucite rod to prevent the products, resulting from the detonation of the tetryl pellets, from obscuring the view of the camera. Above this shield were placed the two tetryl pellets and the detonator. The ionization probe used to trigger the camera and the oscilloscopes was placed at the tetryl-Lucite interface.

To record the reaction two oscilloscopes, a Tektronic No. 535 and a raster oscilloscope were used in conjunction with the smear camera. A spark was arranged to go off at the end of the reaction to provide a common point, on both the oscilloscope and the camera records, from which the time intervals could be measured and compared. The illumination for the camera was obtained from an exploding wire set behind the Lucite rod. Four experiments were performed, two using four-inch long Lucite rods and two using three-inch long Lucite rods.

III. RESULTS

Figure V is a typical record of the attenuation of a shock wave measured by the pressure probes in a Lucite rod using the sweep oscilloscope. The time scale is l μsec per division, and can be read to $^{\pm}$ 0.5 μsec . The alternate positive and negative response of the pressure probes, as they were activated, made it possible to determine the position of any malfunctioning probe. Table I contains the results of the experiments performed



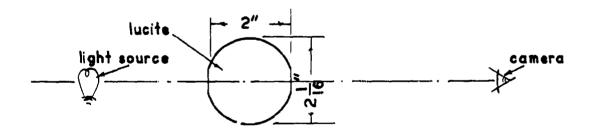
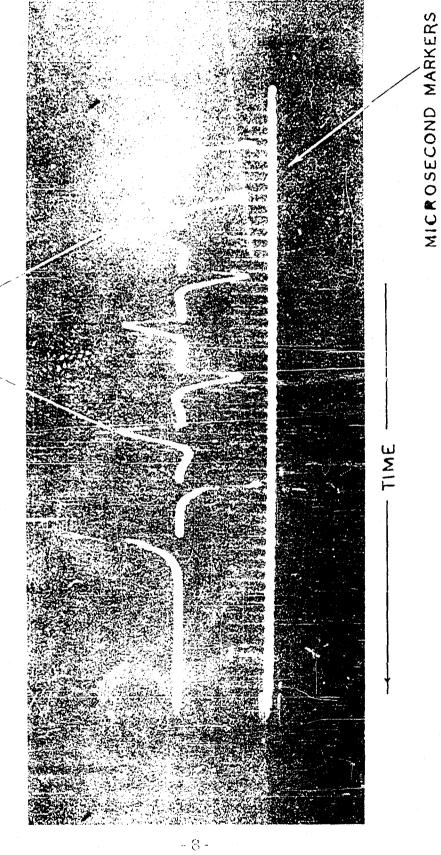


FIGURE IV. - SHOCK ATTENUATION IN LUCITE AND WATER



PRESSURE PROBE RESPONSES

PROBE RESPONSE PRESSURE OF THE RECORD FIGURE V

TABLE I

ATTENUATION OF SHOCK IN A LUCITE ROD (Pressure Probe)

(Two Tetryl Pellets)

	ance of robe		Mean			
in.	mm.	Expt.#1	Time (Micro	Expt.#3	Expt.#4	microsec
0.5	12.7	-	-	2.5	3.0	2.8
1.0	25.4	5.2	6.0	5.5	6.2	5.7
1.5	38.1	8.6	9.6	9.2	9.5	9.2
2.0	50.8	12.0	13.5	12.8	-	12.8
2.5	63.5	16.1	17.8	17.0	17.4	17.1
3.0	76.2	20.5	-	21.1	21.4	21.0
3.5	88.9	25.1	26.5	25.2	26.0	25.7
4.0	101.6	29.7	31.0	29.6	31.1	30.4
4.5	114.3	33.7	37.0	33.7	35.5	35.0
5.0	127.0	-	-	-	-	-
P.078 NO. 1 CO.		ATTENTUATI	on of shock	IN GAP UN	IITS	
0.53	13.4	2.7	-	75 - 0.0	ol in. ace	tate cards
1.00	25.4	6.1	6.0		Lucite dia	
1.50	38.2	9.5	9•5		cite disc	

using two tetryl pellets with the Lucite rods and the gap card units. Table II contains the results of the experiments made with one, three and four tetryl pellets, respectively, and Lucite rods.

Of the four experiments (Expt. #5,6,7 and 8) made using the electronic system and the smear camera, only two could be used for comparison between the two systems. In experiment #7 the fiducial point was not obtained while in experiment #8 the electronic system did not respond satisfactorily. Figure VI shows the records obtained from the smear camera and the raster oscilloscope.

The time scale on the raster oscilloscope was 0.1 microseconds per division and could be read to ± 0.02 microseconds. The time scale for the photographic records was 1.263 mm per microsecond and could be read with a microcomparator to better than ± 0.02 microsecond. The magnification factor for the camera was determined for each experiment by measuring the distance between the probes on the film and relating this to the actual distance between probes. The same magnification factor was used to interpret distance for the shock wave in the water.

Tables III and IV contain the results obtained by the smear camera. The results listed in Table III were obtained by choosing an arbitrary point on the film strip as zero and measuring both time and distance from this point. The data in Table IV were measured from the fiducial point (spark).

TABLE II
ATTENUATION OF SHOCK IN LUCITE ROD (Varying Load)

Dis	tance	Time (mm/microseconds)					
inches	mm.	1-Tetryl	3-Tetry!	7 70 623			
0.5 1.5 1.5 2.5 3.5 4.5 5	12.7 25.4 38.5 50.5 76.9 101.3 127.0	36.4 10.6 10.6 10.6 10.6 10.6 10.6 10.6 10.6	2.6 5.3 8.4 12.3 20.6 25.4 27.4 31.7	2.2 4.6 7.5 11.4 15.8 19.8 24.7 29.4 3			

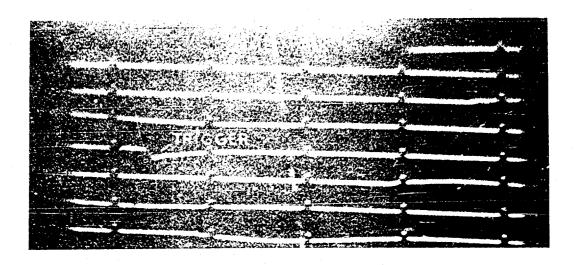


FIGURE VI- A- SECTION OF THE RASTER RECORD

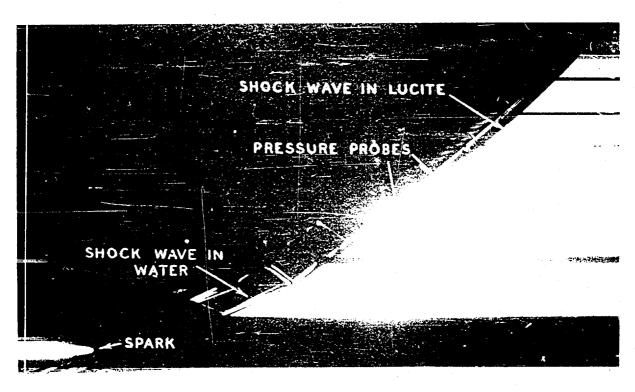


FIGURE VI-B- CAMERA RECORD OF SHOCK WAVE IN LUCITE

AND WATER

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TABLE III
ATTENUATION OF SHOCK IN LUCITE AND WATER*

Shock 1	n	/	Expt.#5,	Lucite	Shock	4	Expt.#6,	4" Lucite
Bilock 7			Time (mm)	Rod Dist.(mm	SHOCK	1n	Time (mm)	Rod Dist.(mm
Luc1te			0 0.536 1.039 1.611 2.446 3.509 4.560 5.556 7.126 8.793 10.289 11.736 13.289	0 0.804 1.316 1.995 2.901 4.005 5.033 7.598 7.598 10.454 11.652 12.956	Lucite		0 0.801 2.282 3.615 4.820 6.299 7.592 9.037 11.100 12.981 14.145 16.223 17.965	0 694 0 694 0 2495 1488 9 9 660 11 9 631 13 9 15
Lucite	in	Н20	15.015 16.746 18.546 20.077 22.157 23.427 25.258 26.517 28.962 30.142	14.334 15.695 17.097 18.813 20.780 22.141 22.991 24.914 25.767 27.410	Lucite	in H ₂ 0	19.317 21.392 23.182 24.955 26.700 28.659 30.284 31.198 32.114 32.897	15.762 17.376 18.687 19.908 21.134 22.534 23.439 24.669 25.197
Interfa Wate	ce	20	35.970 35.156 35.706 36.438 37.406 38.731 40.756 43.471	28.405 29.427 29.659 29.413 30.413 30.993 31.893 33.098	Interfa Wate		33.882 34.972 36.483 38.781 40.911 42.902 44.784 46.406	25.807 26.489 26.489 27.105 28.114 29.867 29.662 31.315

TABLE III, Cont'd.

	Expt.#7,	3" Lucite			Expt.#8,	
Shock in	Rod Time(mm) Dist.(mm)		Shock	1n	Time (mm)	Rod Dist.(mm
Lucite in H ₂ O Interface Water	0.996 9.998 3.366 3.966 7.883 10.975 11.975 15.5280 11.961	1.54.16 1.946 1.946 1.8965 1.8965 1.8966 1.8966 1.8968 1.9986 1.9	Lucite	in H ₂ O	0.574 1.788 1.788 1.788 1.786 1.986 1.986 11.006 11.408 14.422 18.111 19.877 20.463	1124689045015686026572889355 08812508865045686572889355 08813508955602465572889355 11357902456899990123344 11357902248282829333333333333333333333333333333

^{*} The data given here are the readings made directly from the films. The magnification and time factors are given in Table VI.

TABLE IV

RESULTS OF EXPERIMENTS #5 AND #6 USING THE CAMERA RASTER AND SWEEP OSCILLOSCOPE

			Expt. #	5		
Probe No.	Distance from Donor (mm)	Dista Sweep Scope S(µsec)	nce from Raster Scope R(µsec)	Spark Camera C(µsec)	С - S △(µвес)	С - R <u>(</u> µsec)
1 2 3 4 5 6 7 8 5 park	4.2 12.0 12.16 47.4 60.1 785.3	47.43 45.73 45.37 40.59 31.97 28.29 0	47.08 46.65 46.35 40.35 40.35 40.38 28.38 0	48.0 46.31 43.48 43.86 43.86 32.86 28.42 0	+ 0.6 0.5 0.2 0.3 0.7 0.5 1.1	+ 0.32 0.31 0.55 0.00 0.00 1.00
		,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,	Expt. #	6		
1 2 3 4 5 6 7 8 Spark	4.2 12.0 22.1 34.6 47.4 60.1 72.7 85.3	46.52 45.12 42.32 39.25 35.34 26.70 22.50	47.01 45.56 42.77 39.61 35.68 	44.97 42.77 39.61 35.75 31.74 27.71 23.56	- 0.1 + 0.4 0.4 0.4 0.9 1.1	- 0.6 + 0.0 0.1 0.8

IV. DISCUSSION

In the hydrodynamic theory of shock waves, the conservation of momentum requires that

$$P = \beta_0 uU \tag{1}$$

where the initial pressure (P_O) and particle velocity (u_O) are assumed to be zero and where

P = shock pressure

 ρ_0 = initial density of the material

u = particle velocity

U = shock velocity.

In order to obtain the pressure at any point in a shocked homogeneous medium, it is necessary to measure the shock velocity and the particle velocity. However, if a set of data corresponding to equation (1) is known, i.e. the equation of state of the medium is known, a measurement of U vs the attenuation path length (X) for the test geometry can be combined with the known data to give a P-X curve. Since it was desired to use the pressure probes to obtain the U-X data, their adequacy for such measurements was investigated.

A. Pressure Probe Reliability

The construction of the pressure probe (Fig. II) causes a time lag between the arrival of the shock and its recording. The distance between the bare copper wire and the outer copper tube is approximately 0.013 inches. To record the shock, the copper must travel this distance to make contact with the inner core. Moreover, the time lag should increase as the shock pressure and velocity decrease and the response of the pressure probes should fall further behind as the shock is attenuated. In Table IV a comparison is made between data from the smear camera and the sweep oscilloscope (Col. 5) and between the smear camera and the raster oscilloscope (Col. 6). In all but one instance the camera did record the process before the electronic systems did. However, with the exception of probes 7 and 8, placed at a distance of 72.7 and 85.3 mm from the donor, the time lag was, on the whole, less than 0.5 microseconds. The sweep oscilloscope data were slightly higher, 0.6 microseconds. To investigate further the comparative time lags in the systems an equation of the type

$$X = a + bt + ct2 + dt3$$
 (2)

where

- X = distance or gap (mm)
- t = time (microseconds)

was fitted to the data by the electronic computer (IBM 704). In Table V, the first derivatives, $\frac{dX}{dt}$, obtained for experiment #5 are tabulated and compared at various time intervals. A more complete discussion of equation (2) is given in Appendix I.

For the initial 7-10 microseconds, the velocities calculated from the two sets of data differ by only 1%. This interval is the time required for the shock to traverse 1 3/4 inches of Lucite. For 15-20 microseconds $(2\ 1/2-3$ inches of Lucite) the velocities differ by 5-10%. However, the particular equation used did not hold beyond t=20 microseconds and it may be that the difference in velocities at 3 inches or so is somewhat less than indicated.

Thus, the pressure probe may be used to interpret the shock velocity for the initial three inches of Lucite with fair accuracy and reliability (to within 10%). Beyond this, as the shock wave becomes more attenuated, the time lag increases and consequently the percentage error begins to rise very rapidly. Thus, while values at three inches may be off by only 10%, values for larger thicknesses may be in error by much more (e.g., 20% at a four inch thickness). The sensitivity of most propellants and explosives tested, however, lie below the three inch limit. Consequently the pressure probe measurements are considered fairly adequate for this work.

It should be understood that the shock wave velocities, dX/dt, reported and used to compute shock wave pressures in later sections of this report were determined graphically, not analytically. The graphically determined derivatives for both the optical and the probe data are several percent lower than the analytically determined ones of Table V. Moreover, the divergence of the probe results from the optical does not exceed 6% even at four inch thicknesses of Lucite. Indeed, the (dX/dt) vs X curve obtained from Eqn. (2) diverges at both ends of the range 0-20 usec from the graphically determined (dX/dt) vs X; this is the basis for the suggestion above that the inadequacy of the data fit may be responsible for the larger percentage differences of Table V.

B. Velocity vs Distance for Lucite

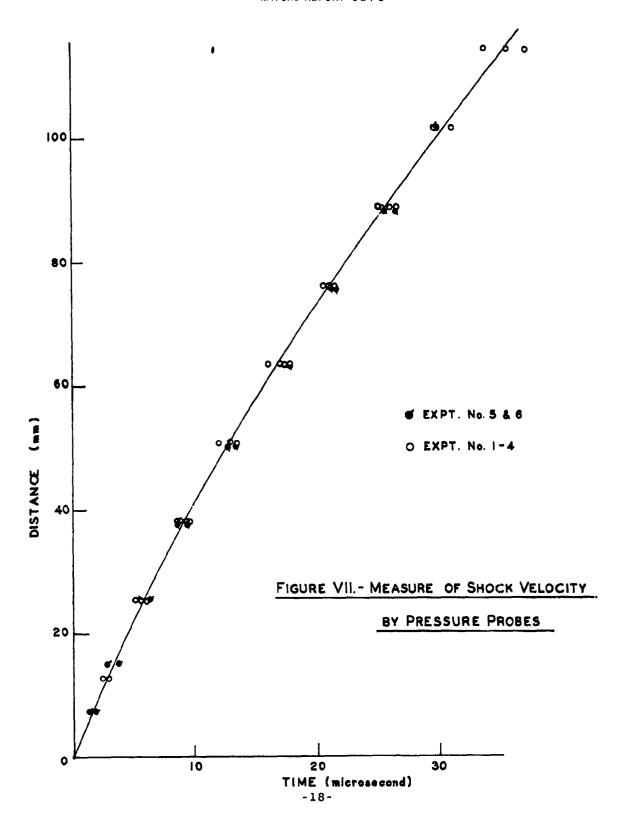
The results of the experiments are plotted in Figures VII, VIII, IX and X. In Figure VII the data obtained with the

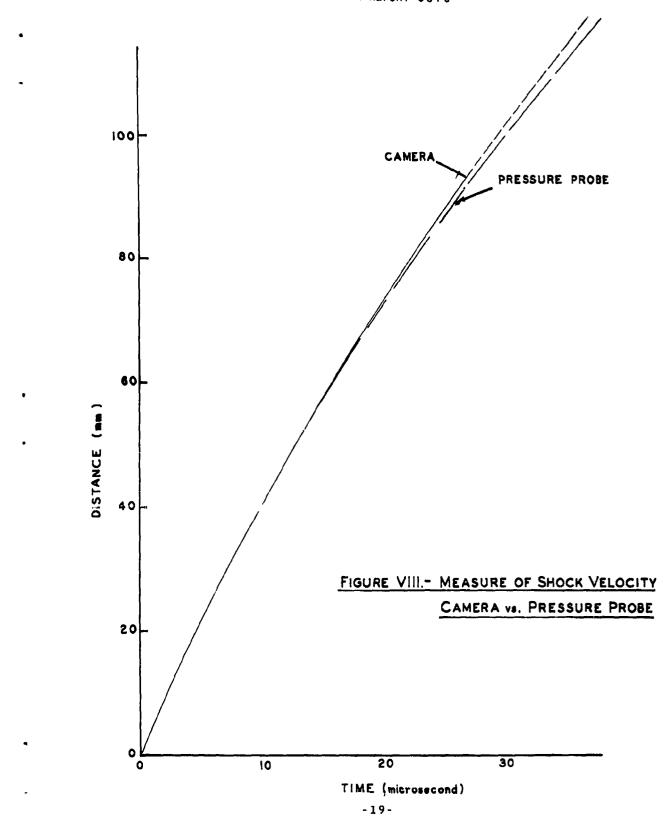
TABLE V

COMPARISON OF THE VELOCITY BETWEEN THE PRESSURE PROBE AND THE SMEAR CAMERA

Time t=microsec.	Pressure Probe dX/dt = mm/microsec.	Camera dX/dt = mm/microsec.	Δ= C - Pp mm/microsec.	% Difference
0	5.0807	5.1233	0.043	0.8
3	4.5277	4.5632	.036	0.8
5	4.2011	4,2398	.039	0.9
7	3.9083	3.9562	.048	1.2
10	3.5322	3,6058	.074	2.0
15	3.0737	3.1216	.148	4.7
50	2.8259	3.0869	.261	8.4

pressure probes are plotted. The precision of these measurements varied from a standard deviation of ± 2.3% to ± 4%. This precision includes any variation due to the probe, the position of the probe, or any variation of the Lucite or the tetryl booster. In Figure VIII a comparison is made between the data obtained with the camera and with the pressure probes. Figure IX compares the measurements made in the gap material with the curve obtained with the camera. It is quite apparent that for the distances measured the cellulose acetate and Lucite systems are comparable. Moreover for these lengths both the pressure probes and the camera give the same results for the same donor. Figure X compares the velocity of the shock fronts obtained from the slopes of the curves in Figure VIII. It is quite apparent that, as the shock was attenuated and the pressure fell, the pressure probe results began to lag behind the results obtained by the smear camera.





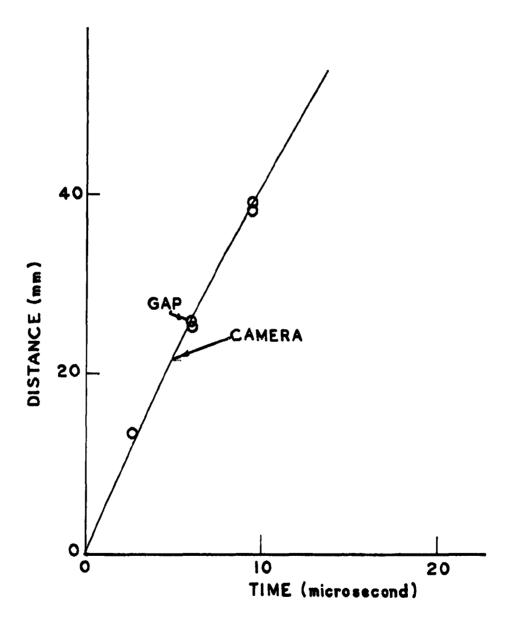
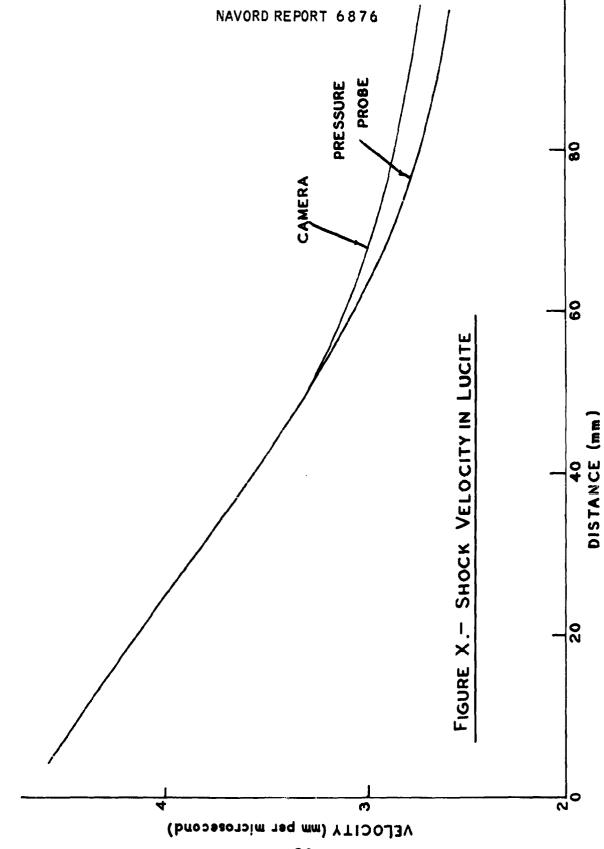


FIGURE IX - COMPARISON - GAP vs CAMERA DATA



In Figure XI the effects of changing the donor load from one to four tetryl pellets is shown.* The detonation pressure and velocity are defined for a given explosive system. If steady state detonation is achieved, these properties (e.g., detonation pressure, and shock velocity induced in Lucite) are not affected by the length of a charge and consequently should not change in going from a donor of one tetryl pellet to two, three, or four pellets. But the transmitted properties in the acceptor (Lucite) do, in fact, depend upon the donor loading (Figure XI). As the number of tetryl pellets is increased, the measured transmitted shock velocity increases. In some cases the transmitted velocity remains constant for a short distance in the Lucite. This was actually measured for the systems of three and four tetryl pellets.

It is known from hydrodynamic theory that the amplitude of the pressure pulse at the interface is a constant and is independent of the length of the donor. It may be assumed that the pressure profile in the explosive adjacent to the interface is fairly constant with the rarefaction following at a finite distance behind the detonation front. The distance between the shock front and the rarefaction should increase with an increase in donor length to an assymptotic value which is achieved at an infinite donor length. Actually it may be possible to attain this value by a practical donor load consisting of five or six tetryl pellets.

The wave transmitted into Lucite is modified in that it will maintain the overall shape of the incident wave, but its amplitude and duration will change. This change is in part due to the impedance mismatch between the donor and acceptor, and the geometry of the system. For short donors the approach and interaction of the rarefaction wave occur after a very short interval of time. The resulting transmitted plateau is of so short a duration that it cannot be detected by the experimental methods used in the present situation. As the donor length is increased (cf. data for 4 pellets) the distance between the shock and rarefaction fronts increases until it is large enough to allow a sufficient length of time for the transmitted plateau to be detected. A further and more comprehensive investigation of these phenomena will be made.

^{*} The preliminary data of Table II indicate a cross-over of some of the curves after a long path of travel through the attenuator. This is attributed to the difficulty in measuring small differences and will be further investigated. Fig. XI has been confined to the area in which the measured differences are large and can be assumed real.

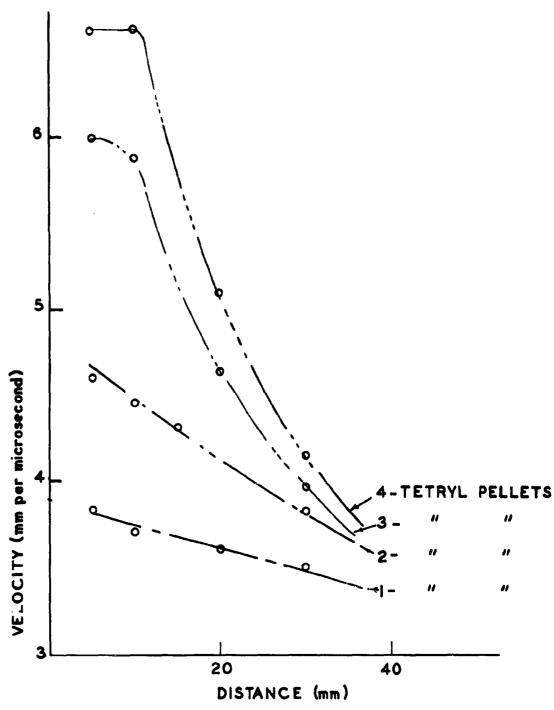


FIGURE XI.- DONOR LOAD AND SHOCK VELOCITY

C. Pressure vs Distance for Lucite

In Figure XII the particle velocity was plotted as a function of the shock velocity from the experimental data obtained on Lucite (5), Plexiglass (Ref. 6 and Appendix II) and Perspex (7). These substances are quite similar in characteristics and it is assumed that their properties in the shock region do not differ from each other. However, the lowest shock strength obtained experimentally is at the upper end of the region critical to this investigation. Most shock sensitivity results on explosives are within the gap range of 30 to 65 mm and the maximum transmitted shock velocity obtained by the two tetryl pellets is about 4.6 mm per microsecond (Figure X). The extrapolation to u = 0 is difficult since the shock pressure is obtained as a product of the particle and shock velocities.

The approximate shock pressures were obtained from the usual boundary approximations (8,9),

$$\mu^{\Gamma} = \mu^{45} O \frac{5(\log n)^{\Gamma}}{(\log n)^{45} O + (\log n)^{\Gamma}}$$
 (3)

where

 μ_{L} = particle velocity in Lucite

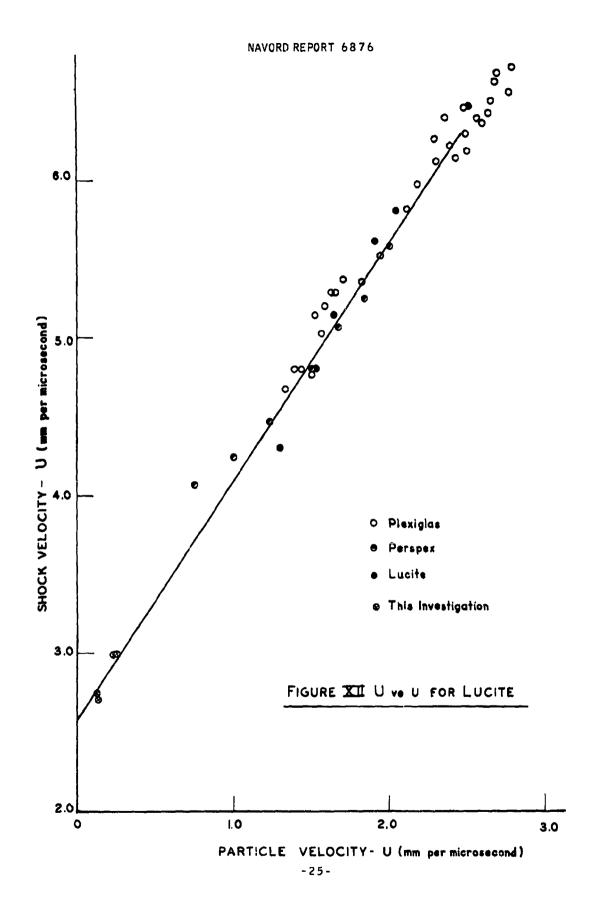
 $\mu_{\rm H_2O}$ = particle velocity in H₂O

% = density of Lucite or water

U = shock velocity in Lucite or water

in conjunction with the experimental data obtained for the shock velocity in Lucite and water, and the particle velocity for water obtained from the literature (10). The calculated particle velocity for Lucite in Eqn. (1) yields the corresponding shock pressure.

Figure XIII is a typical plot of the results (Table IV) obtained by the smear camera. The shock velocities for both Lucite and water are determined at the intersection of the respective curves which corresponds to the Lucite-water interface. Table VI contains the measured shock velocities and the corresponding particle velocities calculated by Eqn. (3). Using these points for the lower pressure region and the other data already available in the higher pressure region a straight line was drawn through all the data. This curve (Fig. XII) was extrapolated to $U=2.59\,\mathrm{mm}$ per microsecond at $\mu=0$; the



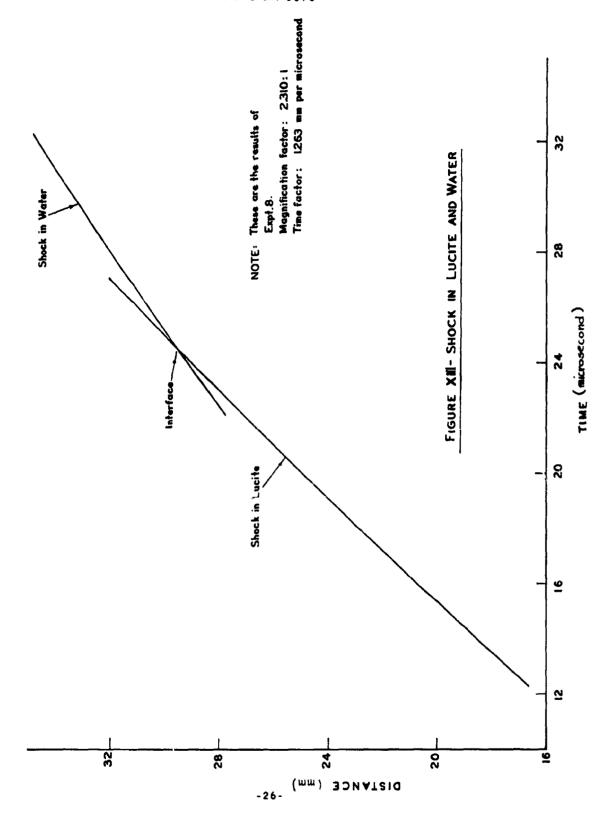


TABLE VI SHOCK VELOCITY IN LUCITE AND HOO, OPTICAL DATA

Expt.No.	Conversion Factor**	UL mm/μsec	UH2O mm/µsec	μ _{H2} O mm/μsec	UL(calc) mm/μsec
5	4.043	2.701	1.840	0.162	0.128
6	4.274	2.744	1.817	0.160	0.125
7	2.938	2.990	2.130	0.312	0.250
8	2.911	2.952	2.069	0.281	0.224

^{**} The conversion factor contains both the magnification factor and time factor.

extrapolated value is approximately equal to the hydrodynamic sound velocity calculated as 2.44 mm per microsecond (Appendix III).

All the data required to develop a pressure-distance curve (P vs X) are available. From experiments 5 thru 8 a U \sim X curve (shock velocity vs distance, Figure X) was obtained for the specified geometry. In addition these experiments provided the data (Table VI) required to calculate and complete the U \sim u curve (shock velocity vs particle velocity, Figure XII). Using these two curves and Eqn. (1) (P \sim \sim \sim \sim 0 to calculate P \sim X (pressure vs distance, Table VII) and obtain the curve in Figure XIV in which the pressure appears to vary exponentially with the distance. Figure XV is a plot of log P vs X and may be approximated by the equation

$$P = 105 e^{-0.0358X}$$
 (4)

These curves will allow direct interpretation of gap length in terms of shock pressure obtained at the end of the Lucite gap. While this pressure is somewhat higher than the pressure entering the acceptor because of the impedance mismatch between the donor and acceptor, it is hoped that this scale of P vs X will offer additional guidance in the sensitivity work. This is especially so since the impedance of Lucite is so near the range found for most explosives. The pressures required to initiate the explosives TNT (34.5 kbar), Composition B (19 Kbar) and tetryl (10 kbar) have been indicated in Figure XIV.

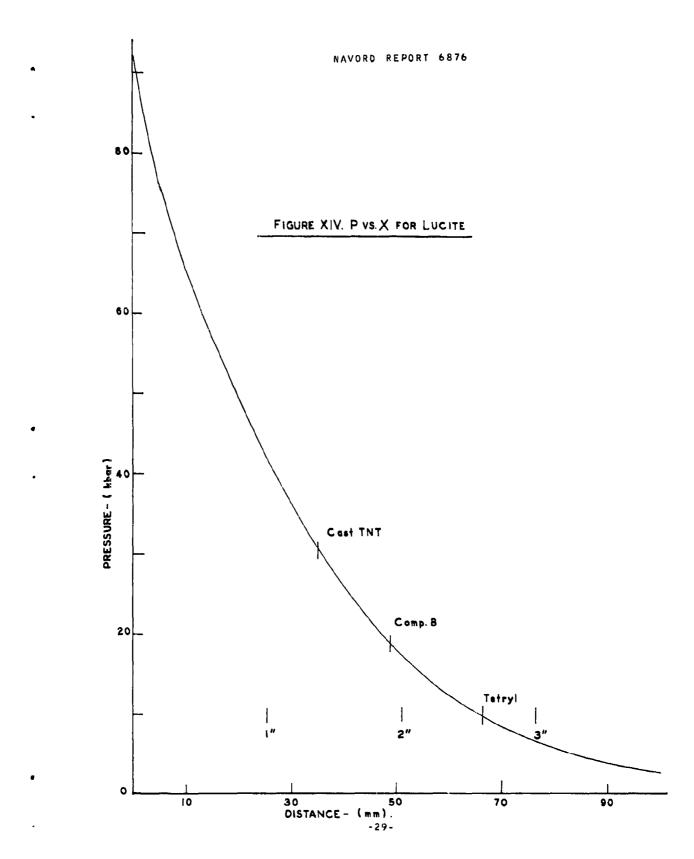
TABLE VII

CALCULATED PRESSURE AND DISTANCE DATA FOR LUCITE

Distance mm	Pressure Kbar	
5 10 20 30 40 50 60 70 80 90 100	75.47 66.08 50.95 36.31 18.24 12.56 4.35 2.89	

ACKNOWLEDGMENT

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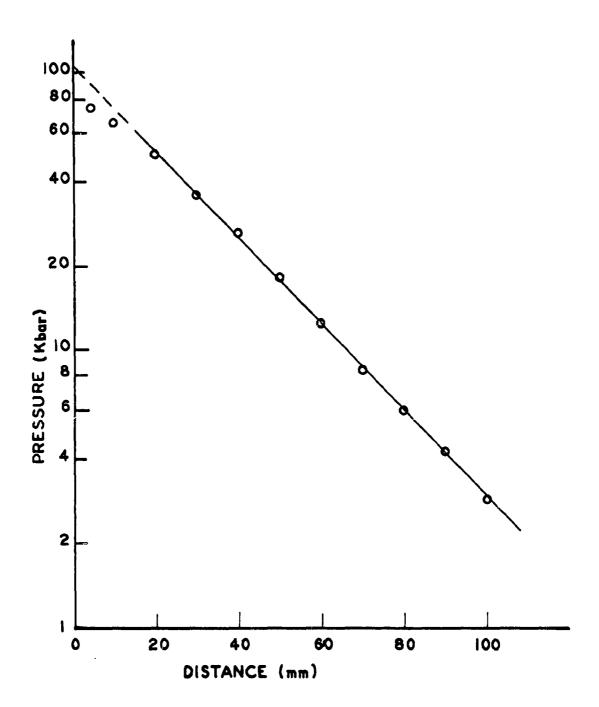


FIGURE XV - LN. PRESSURE vs. DISTANCE

REFERENCES

- 1. A. B. Amster, R. L. Beauregard, G. J. Bryan, and E. K. Lawrence, NavOrd Report 5788, Detonability of Solid Propellants, 1. Test Methods and Instrumentation, 3 February 1958.
- 2. A. B. Amster, R. L. Beauregard, and G. J. Bryan, NavOrd Report 6222, Detonability of Solid Propellants, II. Sensitivity of Some Double Base and Composite Propellants, 15 December 1958, Confidential.
- 3. A. B. Amster, R. L. Beauregard, G. J. Bryan, and E. K. Lawrence, NavOrd Report 6091, Current Status of the Propellant Sensitivity Program at NOL, 20 May 1958, Confidential.
- 4. A. B. Amster and R. L. Beauregard, Rev. Sci. Inst. 30, 942 (1959).
- 5. Los Alamos Scientific Laboratory, private communication.
- 6. N. Coleburn, Naval Ordnance Laboratory, private communication.
- 7. J. S. Buchanan, H. J. James, and G. W. Teague, Armament Research and Development Establishment ARDE Memorandum (MX)20/59, April 1959.
- 8. H. D. Mallory, NavOrd Report 1883, The measurement of Detonation Pressures in Explosives, 5 March 1953, Confidential.
- 9. W. C. Holton, NavOrd Report 3968, The Detonation Pressures in Explosives as Measured by Transmitted Shocks in Water, 1 December 1954, Confidential.
- 10. M. H. Rice and J. M. Walsh, J. Chem. Phys. <u>26</u>, No. 4, 824 (1957).
- C. B. Officer, Introduction to the Theory of Sound Transmission, McGraw-Hill Book Co., Inc., New York, 1958.
- 12. American Institute of Physics Handbook, McGraw-Hill Book Co., Inc., New York, 1957.

APPENDIX I

DERIVED EQUATION FOR x vs t (Distance vs Time)

The least square polynomial equations used to approximate the camera and pressure probe data (Expt. 6) were respectively:

$$X_c = -1.68002 + 5.12334t - 0.100840t^2 + 0.00166433t^3$$
 (1)

$$x_p = -2.84977 + 5.08068t - 0.0984832t^2 + 0.00140377t^3$$
 (2)

The following is a table in which a comparison is made between the distances calculated from the above equation (1) and the distance determined experimentally by the camera.

Time microsec.	X _C Calculated (mm)	X _C Experimental (mm)
1	3.3	4.5
5	21.6	21.7
10	41.1	41.2
20	73.8	74.1
25	89.4	88.6
30	106.2	102.6
35	125.5	115.7

APPENDIX II

EXPERIMENTAL DATA U vs u, - N. L. Coleburn

Below are the experimental data obtained from N. Coleburn determined on Plexiglass ($f_0 = 1.180$).

	7		
Shock Velocity U=m/sec.	Particle Velocity u=m/sec.	Shock Velocity <u>U=m/sec.</u>	Particle Velocity u=m/sec.
6730 6692 6628 6574 6522 6431 6369 6300 6195 6143	2785 2697 26765 2765 2718 2718 25985 2495 2422	4675 4800 4800 4800 4765 5143 5035 5290 5290 5373	1333 1390 1390 1436 1500 1523 1565 1587 1629 1655
6539 6553 6688 6571 6730 6730 6759 6760	2412 2515 2618 2510 2682 2682 2795 2787 2840	5 360 5530 5690 5830 5980 6130 6270 6400	1825 1940 2025 2115 2180 2250 2300 2355 2465
6530 6508 6468 6400 6366 6333 6297 6263 6223	2449 2456 2473 2562 2519 2343 2384 2384	6680	2500

APPENDIX III

CALCULATION OF THE VELOCITY OF SOUND IN LUCITE

The bulk modulus or incompressibility k is defined as the ratio of the hydrostatic pressure on a body to the fractional change in volume.

$$k = \frac{p}{-\Delta}$$

where

k = bulk modulus

p = hydrostatic pressure

Δ = negative dilatation or resultant change in volume

The negative dilatation is propagated with the velocity of

$$\circ - \left(\frac{k}{k}\right)_{1/2}$$

which is the sound velocity.

For an isotropic solid, the bulk modulus may be replaced by the Lamé elastic constants and k is obtained as

$$k - \frac{3\lambda + 2\mu}{3} \tag{3}$$

The sound velocity derived from equations (1) and (3) is

$$c = \left(\frac{3\lambda + 2\mu}{3\rho}\right)^{1/2}$$

A more detailed development of the preceding may be obtained in reference (11).

The Lame constants for Lucite under isothermal conditions were obtained from the American Institute of Physics Handbook (12) as

$$\mu = 0.143 \times 10^{10} \text{ newtons*/m}^2$$

^{*} newton = 105 dynes

 $\lambda = 0.562 \times 10^{10} \text{ newtons/m}^2$

However, the conditions are more nearly adiabati the velocity of sound will be somewhat higher 31

 λ ad = λ iso + Q

where Q is the correction due to heat loss and w 0.044 newtons per m². The velocity of sound in calculated from the above is 2.44 mm per microse

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2nd para. line 15, sentence should read as follows:

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